



# Construction of quantum states by special superpositions of coherent states

G. Mogyorosi<sup>1</sup>, P. Adam<sup>1,2</sup>, E. Molnar<sup>1</sup>, A. Varga<sup>1</sup>, M. Mechler<sup>3</sup>, and J. Janszky<sup>1,2</sup>

<sup>1</sup>Institute of Physics, University of Pécs, H-7624 Pécs, Ifjúság útja 6, Hungary

<sup>2</sup>Institute for Solid State Physics and Optics, Wigner Research Centre for Physics, HAS, H-1525 Budapest, P.O. Box 49, Hungary

<sup>3</sup>MTA-PTE High-Field Terahertz Research Group, H-7624, Pécs, Ifjúság útja 6, Hungary



## Introduction

There has been a persistent interest in the construction and the production of nonclassical states of a harmonic oscillator system due to their potential applications in quantum optics and quantum information processing [1, 2]. Special attention has also been devoted to the idea of quantum state engineering, that is, to the preparation of arbitrary quantum states in the same experimental scheme [3]. An efficient method of quantum state engineering is to construct nonclassical states via discrete coherent-state superpositions. It has been shown that superpositions of even a small number of coherent states placed along a straight line or on a circle in phase space can approximate nonclassical field states with a high degree of accuracy [4, 5, 6]. This protocol exploits the effect of quantum interference between the constituent coherent states.

Inspired by these results, we consider quantum state engineering on an ellipse and on a specific  $3 \times 3$  equidistant lattice of distance  $d$  in phase space. We show, by using an efficient numerical optimization method, that the superpositions of a small number of coherent states of these geometries can approximate certain nonclassical states with a high accuracy. We have optimized numerically the parameters of the chosen geometry and the coefficients of the coherent states via a genetic algorithm [7], in order to obtain the best feasible approximation. In this procedure we determined the misfit parameter [6]

$$\epsilon = 1 - |\langle \psi_N | \Psi \rangle|^2, \quad (1)$$

which is minimized by optimization. The quantity  $|\langle \psi_N | \Psi \rangle|$  is the form of fidelity between the approximating coherent-state superposition  $|\psi_N\rangle$  and the target quantum state  $|\Psi\rangle$ . Obviously, the described method for deriving the approximating coherent-state superpositions can only be used if a well-behaved one-dimensional coherent-state representation of the quantum state exists.

For certain states and parameter ranges the approximation is better than the corresponding one on a circle or along a line.

## Coherent-state superpositions on an ellipse in phase space

Elliptical states can be defined as coherent-state superpositions on an ellipse in phase space

$$|\psi_N\rangle_{\text{ellipse}} = \mathcal{N} \sum_{k=1}^N c_k^{(\text{ellipse})} |\alpha_k\rangle, \quad (2)$$

where

$$\alpha_k = r_k \cdot e^{i\phi_k}, \quad \phi_k = \phi_0 + \frac{2\pi k}{N}, \quad (3)$$

$$r_k = \left( \frac{\cos^2 \phi_k}{a^2} + \frac{\sin^2 \phi_k}{b^2} \right)^{\frac{1}{2}},$$

where  $a$  and  $b$  are arbitrary real numbers. Here the constituent coherent states are placed equidistantly in their phase  $\phi_k$ . Nonclassical properties of such states were analyzed only for states with constant coefficients  $c_k = 1$  [8]. Our task is to find the set of coefficients  $c_k^{(\text{ellipse})}$  in (2) and parameters  $a$ ,  $b$  of the ellipse in (3), which minimize the misfit parameter  $\epsilon^{(\text{ellipse})}$  in (1).

## Coherent-state superpositions on a lattice in phase space

It has been shown that discrete coherent-state superpositions with variable coefficients on a lattice in phase space can be produced in traveling wave optics using only beam splitters and homodyne measurements [9, 10]. We consider the approximation of nonclassical states by such experimentally realizable superpositions on a lattice in phase space.

Let us consider the superposition of 9 coherent states

$$|\psi_9\rangle_{\text{lattice}} = \mathcal{N} \sum_{l=-1}^1 \sum_{k=-1}^1 c_{k,l}^{(\text{lattice})} |l \cdot d + k \cdot id\rangle, \quad (4)$$

on an equidistant lattice centered around the origin in phase space. In this equation  $\mathcal{N}$  is a normalization constant and  $d$  is the distance between adjacent elements of the lattice. As in the case of the ellipse, the task is to find the optimal complex coefficients  $c_{k,l}^{(\text{lattice})}$  and the distance  $d$  in (4). For the measure of the accuracy of the approximation to be optimized, we use the misfit parameter  $\epsilon^{(\text{lattice})}$ .

## References

- [1] Dodonov V V 2002 *J. Opt. B: Quantum Semiclass. Opt.* **4** R1-R33
- [2] Dodonov V V, Man'ko M A, Man'ko V I and Vourdas A 2007 *J. Russ. Laser Res.* **28** 404
- [3] Dell'Anno F, De Siena S and Illuminati F 2006 *Phys. Rep.* **428** 53-168
- [4] Janszky J, Domokos P and Adam P 1993 *Phys. Rev. A* **48** 2213
- [5] Janszky J, Domokos P, Szabó S and Adam P 1995 *Phys. Rev. A* **51** 4191
- [6] Szabo S, Adam P, Janszky J and Domokos P 1996 *Phys. Rev. A* **53** 2698
- [7] Goldberg D. E. 1989 *Genetic Algorithms in Search, Optimization and Machine Learning* (Boston: Addison-Wesley Publishing Co., Inc.)
- [8] Wang Y, Liao Q, Liu Z, Wang J and Liu S 2011 *Opt. Commun.* **284** 282
- [9] Adam P, Kiss T, Mechler M, Darázs Z 2010 *Phys. Scr.* **T140** 014011
- [10] Molnar E, Varga A, Mogyorosi G and Adam P 2014 *21st Central European Workshop on Quantum Optics (Brussels)* (book of abstracts) p 213

## Acknowledgement

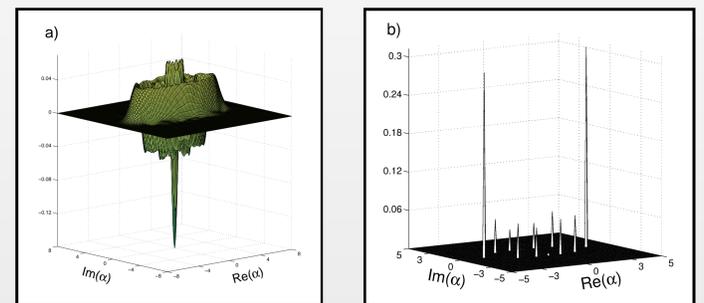
The financial support by SROP-4.1.1.C-12/1/KONV-2012-0005 and by the Hungarian Scientific Research Fund OTKA (Contract No. K83858) is acknowledged.

## Results

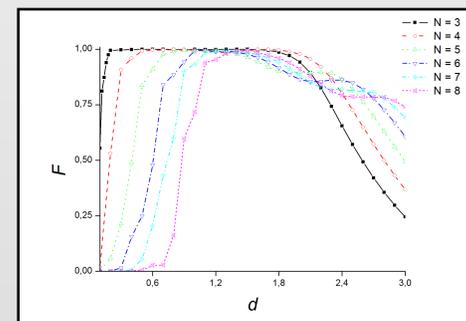
- The minimal misfits  $\epsilon^{(\text{ellipse})}$  and  $\epsilon^{(\text{line})}$  of the approximation of the squeezed number states on an ellipse and along a line with 12 coherent states and the optimal parameters  $a_{\text{opt}}$ ,  $b_{\text{opt}}$  of the ellipse and  $d_{\text{opt}}$  of the line for various photon numbers and for a constant squeezing parameter  $\zeta = 0.5$ .

State	$\epsilon^{(\text{ellipse})}$	$a_{\text{opt}}$	$b_{\text{opt}}$	$\epsilon^{(\text{line})}$	$d_{\text{opt}}$
$ 0, 0.5, 0\rangle$	$6 \times 10^{-4}$	1.25	0.17	$1.4 \times 10^{-5}$	0.27
$ 3, 0.5, 0\rangle$	0.001	2.63	0.97	0.0015	0.47
$ 5, 0.5, 0\rangle$	0.006	3.39	1.44	0.018	0.6
$ 7, 0.5, 0\rangle$	0.0172	4.09	1.54	0.052	0.74

Approximating squeezed number state  $|5, 0.5, 0\rangle$  by  $N = 12$  coherent states on an ellipse (a) Wigner function  $W(\alpha)$  of the approximating coherent-state superposition, (b) pins showing the positions and the absolute values of the coefficients of the constituent coherent states.



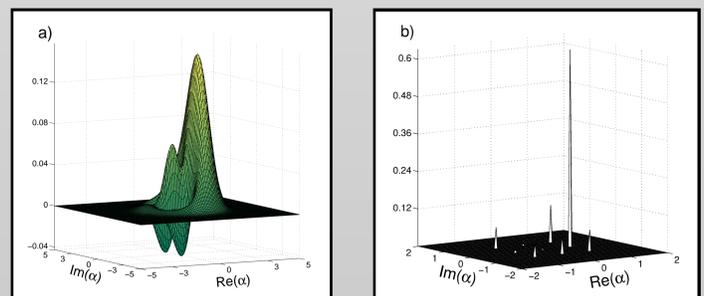
- For the number states we have found that high accuracy can be achieved not only at the optimal value of  $d$  but in a range of distances around this value. The figure shows that as the photon number  $n$  increases, the width of the acceptable range of high accuracy decreases. The lower limit of this range shifts to the direction of larger distances while the upper limit shows a more complicated behavior, but the accuracy of the approximation relevantly decreases for distances  $d > 1.7$ .



- The misfits  $\epsilon^{(\text{lattice})}$  and  $\epsilon^{(\text{circle})}$  of the approximation of number-state superpositions on a lattice and on a circle with 9 coherent states and the optimal distances  $d_{\text{opt}}$  and radii  $R_{\text{opt}}$  for increasing number of superposed states.

State	$\epsilon^{(\text{lattice})}$	$d_{\text{opt}}$	$\epsilon^{(\text{circle})}$	$R_{\text{opt}}$
$\frac{1}{\sqrt{2}}( 0\rangle +  1\rangle)$	$2 \times 10^{-6}$	0.257	$2.2 \times 10^{-5}$	0.54
$\frac{1}{\sqrt{3}}( 0\rangle +  1\rangle +  2\rangle)$	$1.8 \times 10^{-5}$	0.53	$4 \times 10^{-4}$	0.9
$\frac{1}{2}( 0\rangle +  1\rangle +  2\rangle +  3\rangle)$	$6.6 \times 10^{-5}$	0.79	0.0016	1.17
$\frac{1}{\sqrt{5}}( 0\rangle +  1\rangle +  2\rangle +  3\rangle +  4\rangle)$	0.0017	1.37	0.0061	1.32
$\frac{1}{\sqrt{72}}(7 0\rangle + 3 1\rangle + 2 2\rangle +  3\rangle + 3 4\rangle)$	0.0016	1.33	0.0035	1.15

Approximating number state superposition  $\frac{1}{\sqrt{3}}(|0\rangle + |1\rangle + |2\rangle)$  by  $N = 9$  coherent states on a lattice (a) Wigner function  $W(\alpha)$  of the approximating coherentstate superposition, (b) pins showing the positions and the absolute values of the coefficients of the constituent coherent states.



## Conclusion

We have considered the optimal approximation of certain quantum states with the superposition of a finite number of coherent states on an ellipse and on a lattice in phase space. All the coherent-state superpositions we have considered are feasible in current experiments. We have optimized numerically the parameters of the chosen geometry and the coefficients of the coherent states via a genetic algorithm, in order to obtain the best feasible approximation.

First, we placed the coherent states equidistantly in their phase on an ellipse. For squeezed number states in a certain parameter range, the elliptical approximation outperforms the one on a line.

Next we have considered a  $3 \times 3$  equidistant lattice around the origin of phase space. A relatively large set of states, including special numberstate superpositions, appear to be better approximated in this geometry than with a superposition on a circle.