

Construction of quantum states by special superpositions of coherent states

This content has been downloaded from IOPscience. Please scroll down to see the full text.

2015 Phys. Scr. 90 074021

(<http://iopscience.iop.org/1402-4896/90/7/074021>)

View [the table of contents for this issue](#), or go to the [journal homepage](#) for more

Download details:

IP Address: 193.6.62.228

This content was downloaded on 24/06/2015 at 06:25

Please note that [terms and conditions apply](#).

Construction of quantum states by special superpositions of coherent states

P Adam^{1,2}, E Molnar², G Mogyrosi², A Varga², M Mechler³ and J Janszky^{1,2}

¹Institute for Solid State Physics and Optics, Wigner Research Centre for Physics, Hungarian Academy of Sciences, H-1525 Budapest, PO Box 49, Hungary

²Institute of Physics, University of Pécs, H-7624 Pécs, Ifjúság útja 6, Hungary

³MTA-PTE High-Field Terahertz Research Group, H-7624 Pécs, Ifjúság útja 6, Hungary

E-mail: adam.peter@wigner.mta.hu

Received 13 November 2014, revised 14 January 2015

Accepted for publication 1 February 2015

Published 15 June 2015



Abstract

We consider the optimal approximation of certain quantum states of a harmonic oscillator with the superposition of a finite number of coherent states in phase space placed either on an ellipse or on a certain lattice. These scenarios are currently experimentally feasible. The parameters of the ellipse and the lattice and the coefficients of the constituent coherent states are optimized numerically, via a genetic algorithm, in order to obtain the best approximation. It is found that for certain quantum states the obtained approximation is better than the ones known from the literature thus far.

Keywords: coherent-state superpositions, quantum state engineering, nonclassical states

1. Introduction

There has been a persistent interest in the construction and the production of nonclassical states of a harmonic oscillator system due to their potential applications in quantum optics and quantum information processing [1, 2]. Special attention has been devoted to the idea of quantum state engineering, that is, to the preparation of arbitrary quantum states in the same experimental scheme [3]. Several methods have been developed for generating arbitrary superpositions of photon number states of the electromagnetic field in a cavity [4–7]. Experimental schemes have been proposed for the conditional preparation of arbitrary finite superposition of Fock states in a single mode of a traveling-wave optical field by performing alternately coherent quantum-state displacement and photon addition [8, 9] or subtraction [10].

Another efficient method of quantum state engineering is to construct nonclassical states via discrete coherent-state superpositions. It has been shown that superpositions of even a small number of coherent states placed along a straight line or on a circle in phase space can approximate nonclassical field states with a high degree of accuracy [11–13]. This protocol exploits the effect of quantum interference between the constituent coherent states. This phenomenon yields nonclassical properties of the superposition of two coherent states. Such superpositions were originally introduced in [14]

as even and odd coherent states. In the literature these states are generally referred to as Schrödinger cat states due to their correspondence to Schrödinger's famous cat paradox. Properties of such states have been widely discussed in the literature [1, 2, 15–18], and experimental schemes have been developed for their production in different physical systems [19–35]. Several methods have also been proposed for generating discrete coherent-state superpositions on a circle or along a line for electromagnetic field in cavities [13, 36–42] and for the center of mass motion of a trapped ion [43, 44].

A method for producing coherent-state superpositions in traveling-wave optics has been recently proposed [45]. The experimental scheme contains only beam splitters and homodyne detectors and it is capable of producing coherent-state superpositions along the real axis and on an equidistant lattice in phase space. This scheme is based on an apparatus developed previously for generating Schrödinger cat states [28, 29]. The construction of squeezed states and number states by coherent-state superpositions on truncated von Neumann lattices was also considered in [46, 47]. From the results of these papers one can conclude that for the considered states the accuracy of the approximation achievable using the latter lattices is below that of the superposition of the same number of coherent states placed on a circle or along a line, even after optimization by the method of [46]. An experimental scheme for producing coherent-state

superpositions on a truncated von Neumann lattice has been developed for center of mass motion of a trapped ion [48]. Elliptical states have been introduced, too, which are coherent-state superpositions on an ellipse in phase space [49]. An experimental scheme was proposed in the latter reference for generating elliptical states with constant coefficients for the motion of the center of mass of a trapped ion. Nonclassical properties of such states were also analyzed.

Inspired by these results, now we consider quantum state engineering on an ellipse and on a specific lattice in phase space. We show, by using an efficient numerical optimization method, that the superpositions of a small number of coherent states in these geometries can approximate certain nonclassical states with a high accuracy. For certain states and parameter ranges the approximation is better than the corresponding one on a circle or along a line.

The paper is organized as follows. In section 2 we summarize the results on approximation of nonclassical states with coherent-state superpositions on a circle or along a straight line in phase space. In section 3 the construction of squeezed number states by coherent-state superpositions on an ellipse is considered. We present our results on approximating nonclassical states by coherent-state superpositions on a lattice in section 4. Conclusions are given in section 5.

2. Coherent-state superpositions on a circle and along a straight line

In this section we briefly summarize the results known from the literature related to the construction of nonclassical states by coherent-state superpositions along a straight line and on a circle in phase space. A systematic method was developed in [13] for obtaining optimized superpositions from the one-dimensional representation [50–56] of the desired state. In this quantum state engineering protocol a given quantum state $|\Psi\rangle$ is approximated either by an equidistant superposition

$$|\psi_N\rangle_{\text{line}} = \mathcal{N} \sum_{k=1}^N c_k^{(\text{line})} |x_k\rangle \quad (1)$$

of coherent states $|x_k\rangle$ distributed at distances d along the real axis of phase space, or by a superposition

$$|\psi_N\rangle_{\text{circle}} = \mathcal{N} \sum_{k=1}^N c_k^{(\text{circle})} |Re^{i\phi_k}\rangle \quad (2)$$

of coherent states $|Re^{i\phi_k}\rangle$ equally distributed on an arc of a circle in phase space with radius R . In (1) and (2) $\mathcal{N}$ is a normalization constant.

The positions of the coherent states in these superpositions are defined by the following equations:

$$x_k = x_0 + \left(k - \frac{N+1}{2}\right)d, \quad k = 1, \dots, N, \quad (3)$$

$$\phi_k = \phi_0 + \left(k - \frac{N+1}{2}\right)\Delta\phi, \quad k = 1, \dots, N, \quad (4)$$

where x_0 and ϕ_0 determine the center of the distribution, $\Delta\phi$ is the phase difference between two adjacent coherent states.

In the superpositions in (1) and (2), the coefficients c_k are derived from the continuous distribution functions of the one-dimensional coherent-state representations of the state $|\Psi\rangle$ along the real axis

$$|\Psi\rangle = \int_{-\infty}^{\infty} F(x)|x\rangle dx, \quad (5)$$

and on a circle

$$|\Psi\rangle = \int_0^{2\pi} F_R(\phi)|Re^{i\phi}\rangle d\phi, \quad (6)$$

respectively, by the following expressions

$$c_k^{(\text{line})} = F(x_k), \quad (7)$$

$$c_k^{(\text{circle})} = F_R(\phi_k). \quad (8)$$

The optimal values of the distances d and phase differences $\Delta\phi$ can be determined numerically by minimizing the misfit parameter

$$\epsilon = 1 - |\langle\psi_N|\Psi\rangle|^2. \quad (9)$$

The quantity $|\langle\psi_N|\Psi\rangle|^2$ is the fidelity between the approximating coherent-state superposition $|\psi_N\rangle$ and the target quantum state $|\Psi\rangle$.

Obviously, the described method for deriving the approximating coherent-state superpositions can only be used if a well-behaved one-dimensional coherent-state representation of the quantum state $|\Psi\rangle$ exists. For example, the condition of the existence of an appropriate weight function $F_R(\phi)$ on a circle with radius R reads [50]

$$\sum_n \frac{\sqrt{n!}}{R^{n+1}} |c_n| < \infty, \quad (10)$$

where the numbers c_n are the coefficients in the photon number expansion

$$|\Psi\rangle = \sum_{n=0}^{\infty} c_n |n\rangle. \quad (11)$$

We note that for a state with an infinite number of Fock-state coefficients not fulfilling (10) an approximate weight function $F_R^{\text{appr}}(\phi)$ can be derived by appropriate truncation of the expansion of (11). As a consequence, discrete coherent-state superpositions approximating a nonclassical state can always be constructed on a circle in phase space.

Using this procedure, approximating discrete coherent-state superpositions were determined in [13] for various nonclassical states, including displaced squeezed number states, squeezed coherent states, binomial states, Hermite-polynomial states, and amplitude squeezed states. In [57] coherent-state superpositions approximating phase-optimized states on the real axis of phase space were derived by a different numerical method. This method varies not only the coefficients but also the positions of the coherent states in phase space. Though there is no one-dimensional straight-line

coherent-state representation known for these states, the numerically obtained superpositions of high accuracy consist of a few coherent states distributed nearly equidistantly. This experience can be explained by the properties of the quantum interference between the constituent coherent states. The interference phenomenon arises typically at a certain range of distance between coherent states. In view of this, the choice of the equidistant states in the superpositions (1) and (2) seems to be reasonable. We note that the accuracy of the approximation can be generally increased by increasing the number of constituent states in the superpositions.

From these results one can conclude that any state can be approximated by discrete coherent-state superpositions. Even expecting a given accuracy, different superpositions of various geometries can approximate the same nonclassical state. In general, small number of constituent states are required if the geometry of their positions fits well to the state to be approximated. For example, states exhibiting circular symmetry in the Wigner function such as photon number states or amplitude squeezed states can be effectively approximated by superpositions on a circle. The optimal value of the radius of the circle fits to the value of the photon number [13] when the constituent coherent states in the superposition are located under the central part of the Wigner function of the target state. On the other hand, effective approximation of quadrature squeezed states can be realized by superpositions along a line. These facts also motivate our considerations of coherent-state superpositions on an ellipse for approximating squeezed number states in the next section.

As we have already mentioned in the introduction, several schemes have been proposed in the literature for producing coherent-state superpositions approximating quantum states along a line and on a circle in cavities [13, 36–42] and trapped ion systems [43, 44]. The experimental realization of these quantum state engineering protocols seems easier than that of the protocols based on the realization of the number-state expansions [4–7] of the target states. The reason is that the complexity of the experiment, for example, the number of atoms sent through the cavity or the number of the applied laser pulses, is proportional to the number of terms in the number-state expansion of the state [5] or the number of coherent states in the approximating superposition [13, 43, 44]. According to the results we have summarized here thus far, for certain quantum states at a given accuracy, the number of the required coherent states of the approximation can be less than that of the number states.

3. Coherent-state superpositions on an ellipse

Elliptical states has been introduced in [49]. An experimental scheme is also proposed in this reference for generating equidistant coherent-state superpositions with constant coefficients on an ellipse for the motion of the center of mass of a trapped ion. This scheme can probably be generalized for producing arbitrary superpositions.

Elliptical states can be defined as coherent-state superpositions on an ellipse in phase space

$$|\psi_N\rangle_{\text{ellipse}} = \mathcal{N} \sum_{k=1}^N c_k^{(\text{ellipse})} |\alpha_k\rangle, \quad (12)$$

where

$$\alpha_k = r_k \exp(i\phi_k), \quad (13)$$

$$\phi_k = \phi_0 + \frac{2\pi k}{N}, \quad (14)$$

$$r_k = \left(\frac{\cos^2 \phi_k}{a^2} + \frac{\sin^2 \phi_k}{b^2} \right)^{-1/2}, \quad (15)$$

where a and b are arbitrary real numbers. Here the constituent coherent states are placed equidistantly in their phase ϕ_k . Nonclassical properties of such states were analyzed only for states with constant coefficients $c_k = 1$ [49].

Here we consider the construction of quantum states via elliptical states. Inspired by the considerations in the previous section, we choose displaced squeezed number states for demonstrating the idea of approximating quantum states by discrete coherent-state superpositions on an ellipse.

Displaced squeezed states are defined as

$$|n, \zeta, Z\rangle = \hat{D}(Z) \hat{S}(\zeta) |n\rangle, \quad (16)$$

where $\hat{D}(Z)$ is the displacement operator

$$\hat{D}(Z) = \exp(Z\hat{a}^\dagger - Z^*\hat{a}), \quad (17)$$

while $\hat{S}(\zeta)$ is the squeezing operator

$$\hat{S}(\zeta) = \exp\left(\frac{1}{2}\zeta^*\hat{a}^2 - \frac{1}{2}\zeta\hat{a}^{\dagger 2}\right), \quad \zeta = re^{i\theta}. \quad (18)$$

In the above definitions $\hat{a}$ and $\hat{a}^\dagger$ are the annihilation and creation operators, respectively, Z is the amount of displacement in phase space, and ζ is the complex squeezing parameter.

The Wigner function for a squeezed number state reads

$$W_{\text{SN}}(\alpha) = \frac{2}{\pi} \exp\left[\frac{1}{2}(\alpha - \alpha^*)^2 e^{-2r} - \frac{1}{2}(\alpha + \alpha^*)^2 e^{2r}\right] \times (-1)^n \mathcal{L}_n\left[(\alpha + \alpha^*)^2 e^{2r} - (\alpha - \alpha^*)^2 e^{-2r}\right], \quad (19)$$

where $\mathcal{L}_n(x)$ is the Laguerre polynomial of degree n :

$$\mathcal{L}_n(x) = \sum_{m=0}^n (-1)^m \binom{n}{m} \frac{x^m}{m!}. \quad (20)$$

In figure 1 the Wigner function $W(\alpha)$ of the state $|3, 0.5, 0\rangle$ is presented. It can be seen that the Wigner function of the squeezed number states has an elliptic shape. We note that for the displaced number state the one-dimensional coherent-state distribution function is also known, so approximating discrete coherent-state superpositions along the real axis can be easily obtained using the method described in section 2.

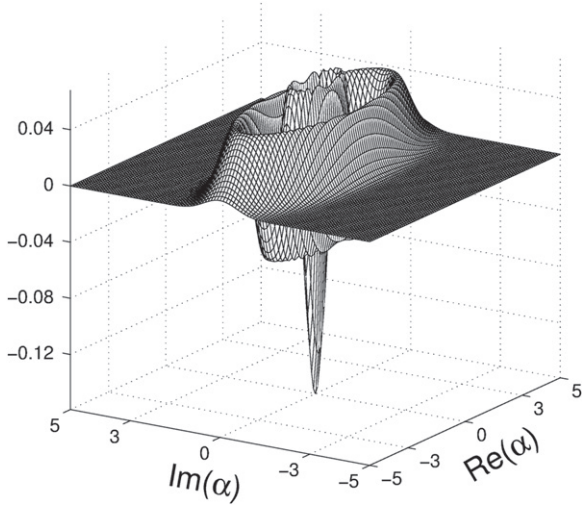


Figure 1. Wigner function $W_{\text{SN}}(\alpha)$ for the squeezed number state $|3, 0.5, 0\rangle$.

Table 1. The minimal misfits $\epsilon^{(\text{ellipse})}$ and $\epsilon^{(\text{line})}$ of the approximation of the squeezed number states on an ellipse and along a line with 12 coherent states and the optimal parameters a_{opt} , b_{opt} of the ellipse and d_{opt} of the line for various photon numbers and for a constant squeezing parameter $\zeta = 0.5$.

State	$\epsilon^{(\text{ellipse})}$	a_{opt}	b_{opt}	$\epsilon^{(\text{line})}$	d_{opt}
$ 0, 0.5, 0\rangle$	6×10^{-4}	1.26	0.17	1.4×10^{-5}	0.27
$ 3, 0.5, 0\rangle$	0.001	2.63	0.97	0.0015	0.47
$ 5, 0.5, 0\rangle$	0.006	3.39	1.44	0.018	0.6
$ 7, 0.5, 0\rangle$	0.0172	4.09	1.54	0.052	0.74

Our task is to find the set of coefficients $c_k^{(\text{ellipse})}$ in (12) and parameters a, b of the ellipse in (15), which minimize the misfit parameter in (9). Since it is not a convex function of the parameters, an algorithm supporting non-convex optimization has to be chosen. In our calculations we have chosen to use a genetic algorithm [58] for solving the optimization problem, as its applicability is quite general and we found it effective in the present case. In order to eliminate the uncertainty of the method due to its stochastic nature, we have performed multiple runs and we chose the parameter set for which the value of the misfit parameter (9) was minimal. Besides, at a given accuracy, various solutions with slightly different parameters can be found for a given target state. This degeneracy of the optimum can be foreseen from the general properties of the discrete coherent-state superposition described in the last paragraph of section 2.

In tables 1–4 we present our results on approximating displaced squeezed number states on an ellipse with 12 coherent states for various values of the photon number, the squeezing parameter and the amount of displacement. We show the misfit $\epsilon^{(\text{ellipse})}$ defined in (9) and the ellipse parameters a_{opt} and b_{opt} introduced in (15) corresponding to the optimum. For comparison, the tables also show the misfit $\epsilon^{(\text{line})}$ and the optimal distance d_{opt} for approximating displaced squeezed number states along the real axis of phase

Table 2. The misfits $\epsilon^{(\text{ellipse})}$ and $\epsilon^{(\text{line})}$ of the approximation of the squeezed number states on an ellipse and along a line with 12 coherent states and the optimal parameters a_{opt} , b_{opt} of the ellipse and d_{opt} of the line for various squeezing parameters and for a constant photon number $n = 3$.

State	$\epsilon^{(\text{ellipse})}$	a_{opt}	b_{opt}	$\epsilon^{(\text{line})}$	d_{opt}
$ 3, 0.1, 0\rangle$	5×10^{-5}	1.48	1.56	0.0099	0.3
$ 3, 0.3, 0\rangle$	2.7×10^{-4}	2.07	0.89	0.0035	0.41
$ 3, 0.5, 0\rangle$	0.001	2.63	0.97	0.0015	0.53
$ 3, 0.7, 0\rangle$	0.011	3.12	0.94	0.0033	0.71
$ 3, 0.8, 0\rangle$	0.017	3.55	0.99	0.0028	0.76
$ 3, 1.2, 0\rangle$	0.12	4.91	0.75	0.0086	1.16

Table 3. The misfits $\epsilon^{(\text{ellipse})}$ and $\epsilon^{(\text{line})}$ of the approximation of the displaced squeezed number states on an ellipse and along a line with 12 coherent states and the optimal parameters a_{opt} , b_{opt} of the ellipse and d_{opt} of the line for a constant squeezing parameter $\zeta = 0.5$, a constant photon number $n = 3$, and for various displacements Z .

State	$\epsilon^{(\text{ellipse})}$	a_{opt}	b_{opt}	$\epsilon^{(\text{line})}$	d_{opt}
$ 3, 0.5, 0.2i\rangle$	0.0041	2.62	1	0.0078	0.46
$ 3, 0.5, 0.3i\rangle$	0.0035	2.64	1.33	0.0046	0.67
$ 3, 0.5, 0.5i\rangle$	0.009	2.49	1.6	0.028	0.7
$ 3, 0.5, 0.8i\rangle$	0.038	2.37	1.73	0.12	0.62
$ 3, 0.5, i\rangle$	0.088	2.58	1.72	0.29	0.53
$ 3, 0.5, 1.5i\rangle$	0.026	2.87	1.94	0.47	0.5
$ 3, 0.5, 1.8i\rangle$	0.068	2.63	2.52	0.56	0.55
$ 3, 0.5, 2i\rangle$	0.12	2.5	3.41	0.65	0.24

Table 4. The misfits $\epsilon^{(\text{ellipse})}$ and $\epsilon^{(\text{line})}$ of the approximation of the squeezed number states on an ellipse and along a line with 12 coherent states and the optimal parameters a_{opt} , b_{opt} of the ellipse and d_{opt} of the line for various phase θ and a constant magnitude $r = 0.5$ of squeezing and a constant photon number $n = 3$.

State	$\epsilon^{(\text{ellipse})}$	a_{opt}	b_{opt}	$\epsilon^{(\text{line})}$	d_{opt}
$ 3, 0.5, 0\rangle$	0.001	2.63	0.97	0.0015	0.53
$ 3, 0.5 \exp(i\pi/12), 0\rangle$	0.006	2.77	1.33	0.0018	0.52
$ 3, 0.5 \exp(i\pi/6), 0\rangle$	0.017	2.62	1.61	0.012	0.47
$ 3, 0.5 \exp(i\pi/4), 0\rangle$	0.03	2.37	1.6	0.09	0.46
$ 3, 0.5 \exp(i\pi/2), 0\rangle$	0.022	2.06	2.01	0.29	0.4
$ 3, 0.5 \exp(i5\pi/6), 0\rangle$	0.0173	1.64	2.6	0.59	0.29
$ 3, 0.5 \exp(i11\pi/12), 0\rangle$	0.0067	1.42	2.66	0.55	0.24
$ 3, 0.5 \exp(i\pi), 0\rangle$	0.0011	0.97	2.59	0.6	0.28

space with 12 equidistant coherent states. We note that we determined the coefficients $c_k^{(\text{line})}$ and the optimal distances d_{opt} of the constituent coherent states in the line superpositions by using the same genetic algorithm. This method may lead to a better approximation than the one described in section 2.

In table 1 our results are presented for a fixed real squeezing parameter $\zeta = 0.5$ and no displacement, while the photon number n is changed. It can be seen that increasing the photon number leads to the growth of the size of the optimal

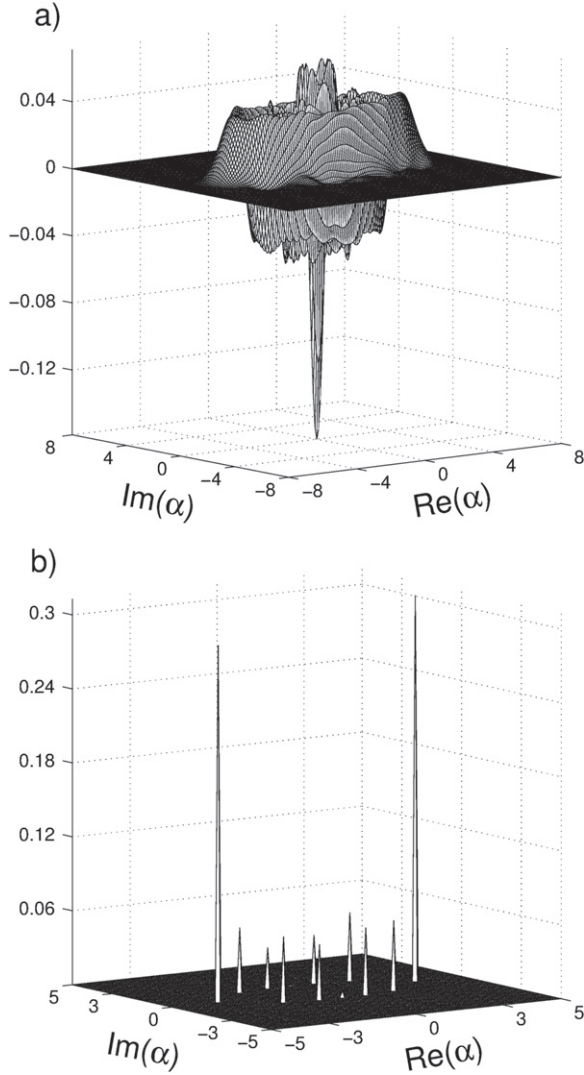


Figure 2. Squeezed number state $[5, 0.5, 0]$ approximated by $N = 12$ coherent states on a ellipse (a) Wigner function $W(\alpha)$ of the approximating coherent-state superposition, (b) pins showing the positions and the absolute values of the coefficients of the constituent coherent states.

ellipse while the accuracy of the approximation ϵ slightly decreases but remains rather high for all the states considered. According to the table the approximation on an ellipse proves to be better than approximation along a line except for the squeezed vacuum state $[0, 0.5, 0]$. This effect confirms the intuition described at the end of the previous section that the best performance can be achieved by fitting the location of the constituent coherent states in the superposition to the geometry and the size of the Wigner function of the target state. Figure 2 shows the pins representing the positions and the absolute values of the coefficients of the $N = 12$ constituent coherent states in the coherent-state superposition approximating the squeezed number state $[5, 0.5, 0]$ on an ellipse. The Wigner function $W(\alpha)$ of the approximating superposition is also shown. This function practically coincides with the Wigner function of the target state, only some imperfections

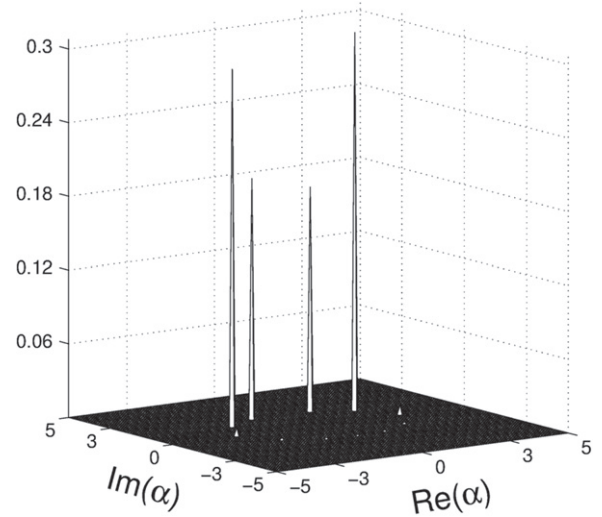


Figure 3. Pins showing the positions and the absolute values of the coefficients of the constituent coherent states for approximating squeezed displaced number state $[3, 0.5, 1.5i]$ with $N = 9$ coherent states on a lattice.

can be seen appearing in the form of small fluctuations due to quantum interference.

Table 2 contains results for constant photon number $n = 3$ without displacement and for various real squeezing parameters ζ . It can be seen from the table that for stronger squeezing the optimal ellipse becomes more and more elongated, that is, the ratio a/b of the parameters of the ellipse increases, in the same way as the Wigner function of the state becomes more elongated. The data show that the superposition on an ellipse proves to be better only for smaller squeezing parameters. For a stronger squeezing when the Wigner function of the state is dominantly elongated along the real axis, superpositions along the appropriate line perform better.

The results for displaced squeezed number states are shown in table 3. In this case both the photon number $n = 3$ and the squeezing parameter $\zeta = 0.5$ are kept constant, while the amount of displacement Z is changed. According to the table, increasing the displacement results generally in a decrease of the accuracy, but exceptions can occur at certain values of the displacement. Though the tendency for approximating along a line is the same, the approximation is much less accurate. The data also show that the optimal ellipse tends to deform towards the displacement ensuring the interfering coherent states to be close to the area of the Wigner function of the displaced squeezed target state. Interestingly, the optimal length a_{opt} of the major axis of the ellipse does not change monotonously. From figure 3 it can be seen that for a larger displacement $Z = 1.5i$ only superposed coherent states in the upper half of phase space have non-negligible weights.

Finally, table 4 presents results on squeezed number states with various phases θ but constant magnitude $r = 0.5$ of squeezing, keeping the photon number $n = 3$ fixed. From the data it can be seen that the deformation of the ellipse changes in accordance with the increase of the phase of squeezing.

The originally minor axis b of the ellipse is growing, for phases larger than $\pi/2$ the ellipse elongates in the direction of the imaginary axis of phase space. The change in accuracy is not monotonous, the accuracy is the highest for squeezing corresponding to the orientation of the ellipse, that is, for $\theta = 0$ and $\theta = \pi$. For larger phases of squeezing ($\theta > \pi/6$) approximation on an ellipse proves to be better than the one along a line.

4. Discrete coherent-state superpositions on a lattice in phase space

Recently, it has been shown that discrete coherent-state superpositions with variable coefficients on a lattice in phase space can be produced in traveling-wave optics using only beam splitters and homodyne measurements [45]. In this section we consider the approximation of nonclassical states by such *experimentally realizable* superpositions on a lattice in phase space.

Accordingly, let us consider the superposition of 9 coherent states

$$|\psi_9\rangle_{\text{lattice}} = \mathcal{N} \sum_{l=-1}^1 \sum_{k=-1}^1 c_{k,l}^{(\text{lattice})} |l \cdot d + k \cdot id\rangle, \quad (21)$$

on an equidistant lattice centered around the origin in phase space. In this equation d is the distance between adjacent elements of the lattice and $\mathcal{N}$ is a normalization constant.

Again, we have used the same genetic algorithm as in the previous section for finding the optimal complex coefficients $c_{k,l}^{(\text{lattice})}$ and the distance d in (21). For the measure of the accuracy of the approximation to be optimized, we use the misfit parameter $\epsilon^{(\text{lattice})}$ introduced in (9). For comparison, we also determined the misfit $\epsilon^{(\text{circle})}$ and the optimal radius R_{opt} for equidistant coherent-state superpositions consisting 9 coherent states on a circle and approximating the given target state.

First we examine the problem of constructing number states by lattice and circle superpositions. The results are presented in table 5. It can be seen that all the considered number states ($n \leq 8$) can be approximated with high accuracy on both formations. The data show that the approximation on a circle is better than the corresponding one on a lattice for number states $|n\rangle$ except for $n = 1$. The optimal distance d_{opt} and radius R_{opt} increases with n , that is, with the growth of the area of the Wigner function of the target state. In figure 4 we present the pins showing the positions and the absolute values of the coefficients of the $N = 9$ constituent coherent states in the coherent-state superposition approximating the number state $|3\rangle$ on a lattice and the Wigner function of the approximating state which is almost identical with the one of the target state.

For the number states we have found that high accuracy can be achieved not only at the optimal value of d but in a range of distances around this value. This finding is shown in figure 5 where the quantity $F = \sqrt{1 - \epsilon^{(\text{lattice})}}$ is presented as a function of n . The quantity F is the absolute value of the

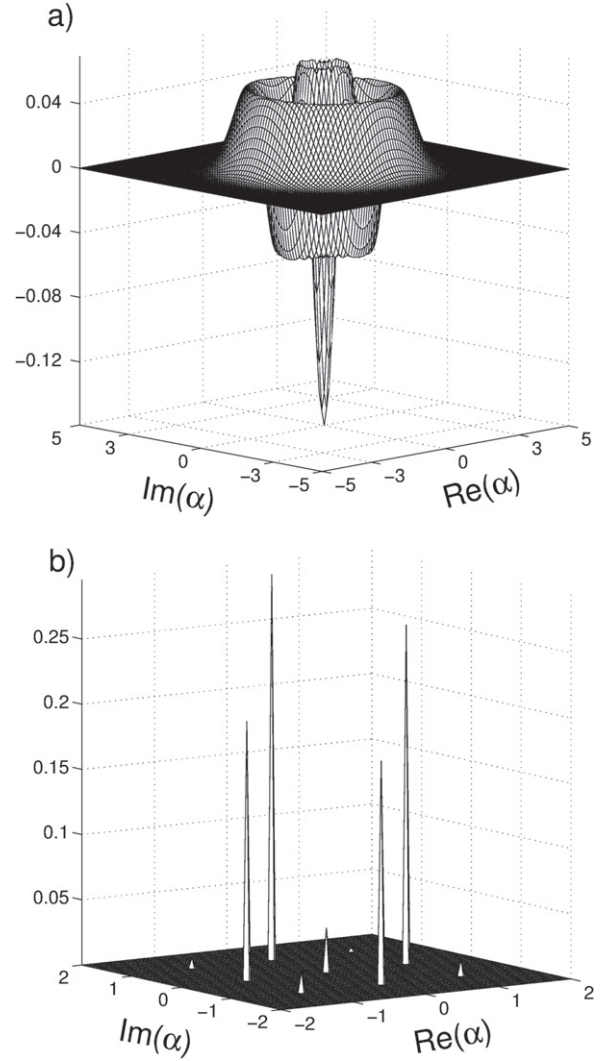


Figure 4. Number state $|3\rangle$ approximated by $N = 9$ coherent states on a lattice (a) Wigner function $W(\alpha)$ of the approximating coherent-state superposition, (b) pins showing the positions and the absolute values of the coefficients of the constituent coherent states.

Table 5. The misfits $\epsilon^{(\text{lattice})}$ and $\epsilon^{(\text{circle})}$ of the approximation of the number states on a lattice and on a circle with 9 coherent states and the optimal distances d_{opt} and radii R_{opt} for various photon numbers n .

State	$\epsilon^{(\text{lattice})}$	d_{opt}	$\epsilon^{(\text{circle})}$	R_{opt}
$ 1\rangle$	2.2×10^{-6}	0.22	3.4×10^{-6}	0.23
$ 2\rangle$	3.1×10^{-4}	0.39	5.6×10^{-6}	0.3
$ 3\rangle$	3.5×10^{-4}	0.5	3.9×10^{-5}	0.55
$ 4\rangle$	5×10^{-4}	0.73	1.3×10^{-5}	0.8
$ 5\rangle$	0.004	1.03	2.8×10^{-4}	1.25
$ 6\rangle$	0.0085	1.2	4.8×10^{-5}	1.73
$ 7\rangle$	0.009	1.35	8.3×10^{-5}	1.92
$ 8\rangle$	0.0097	1.5	2.1×10^{-4}	1.99

scalar product $|\langle \psi_N | \Psi \rangle|$ sometimes used as a form of fidelity in the literature. The figure shows that as the photon number n increases, the width of the acceptable range of high accuracy decreases. The lower limit of this range shifts to the direction

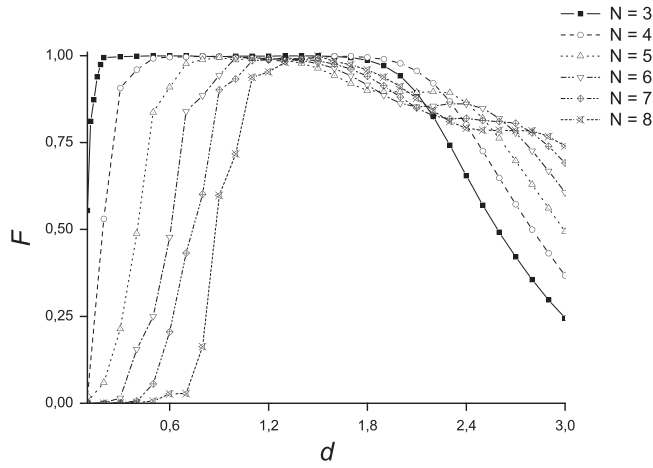


Figure 5. Precisions of approximating number states $|N\rangle$ on a lattice depending on the parameter d .

Table 6. The misfits $\epsilon^{(\text{lattice})}$ and $\epsilon^{(\text{circle})}$ of the approximation of number-state superpositions on a lattice and on a circle with 9 coherent states and the optimal distances d_{opt} and radii R_{opt} for increasing number of superposed states.

State	$\epsilon^{(\text{lattice})}$	d_{opt}	$\epsilon^{(\text{circle})}$	R_{opt}
Ψ_{01}	2×10^{-6}	0.257	2.2×10^{-5}	0.54
Ψ_{012}	1.8×10^{-5}	0.53	4×10^{-4}	0.9
Ψ_{0123}	6.6×10^{-5}	0.79	0.0016	1.17
Ψ_{01234}	0.0017	1.37	0.0061	1.32
$\tilde{\Psi}_{01234}$	0.0016	1.33	0.0035	1.15

of larger distances while the upper limit shows a more complicated behavior, but the accuracy of the approximation relevantly decreases for distances $d > 1.7$.

Let us now consider the approximation of the number-state superpositions

$$\Psi_{01} = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle), \quad (22)$$

$$\Psi_{012} = \frac{1}{\sqrt{3}}(|0\rangle + |1\rangle + |2\rangle), \quad (23)$$

$$\Psi_{0123} = \frac{1}{2}(|0\rangle + |1\rangle + |2\rangle + |3\rangle), \quad (24)$$

$$\Psi_{01234} = \frac{1}{\sqrt{5}}(|0\rangle + |1\rangle + |2\rangle + |3\rangle + |4\rangle), \quad (25)$$

$$\tilde{\Psi}_{01234} = \frac{1}{\sqrt{72}}(7|0\rangle + 3|1\rangle + 2|2\rangle + |3\rangle + 3|4\rangle) \quad (26)$$

by lattice and circle coherent-state superpositions. In table 6 we present the misfits $\epsilon^{(\text{lattice})}$ and $\epsilon^{(\text{circle})}$ of the approximations and the optimal distances d_{opt} of adjacent lattice elements, and the optimal radii R_{opt} . The table shows that the accuracy of the approximation on a lattice is quite high for all the considered superpositions, and it is always higher than the

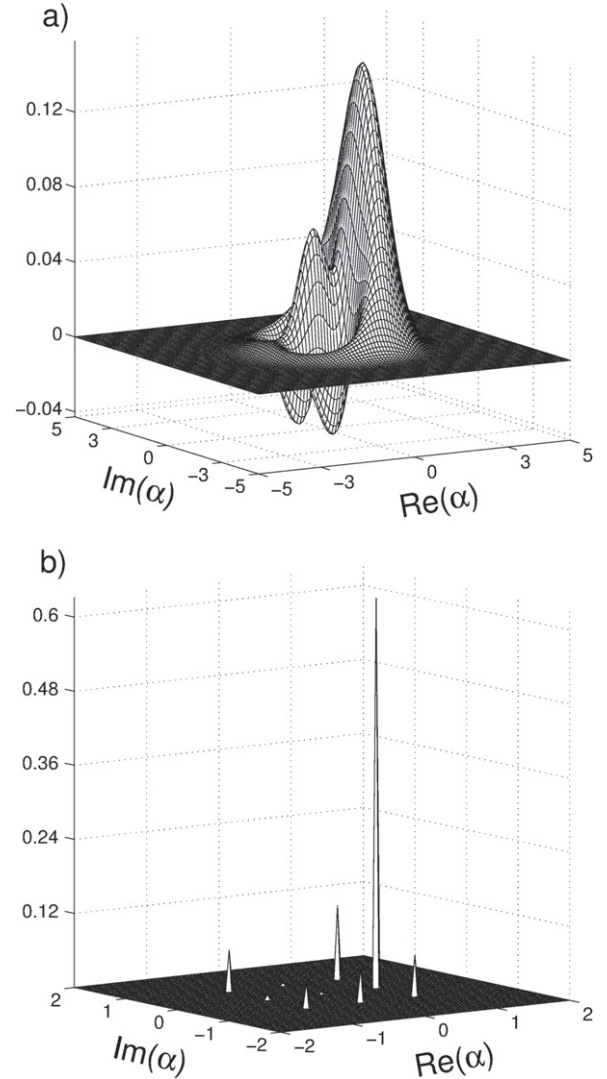


Figure 6. Number-state superposition $\frac{1}{\sqrt{3}}(|0\rangle + |1\rangle + |2\rangle)$ approximated by $N = 9$ coherent states on a lattice (a) Wigner function $W(\alpha)$ of the approximating coherent-state superposition, (b) pins showing the positions and the absolute values of the coefficients of the constituent coherent states.

accuracy of the approximation on a circle. Increasing the number of constituent number states in the target state including higher number states results in larger optimal distances d_{opt} and optimal radii R_{opt} . We note that generation of such types of superpositions can be important for some quantum information protocols [59–61]. Several schemes have been developed for producing such states in the literature [61–63]. From the results presented here it can be seen that the construction and generation of these states by coherent-state superposition on a lattice can be a rather effective method.

In figure 6 we present the pins showing the positions and the absolute values of the coefficients of the $N = 9$ constituent coherent states in the coherent-state superposition approximating the $\frac{1}{\sqrt{3}}(|0\rangle + |1\rangle + |2\rangle)$ state on a lattice and the Wigner function of the approximating state. This function is

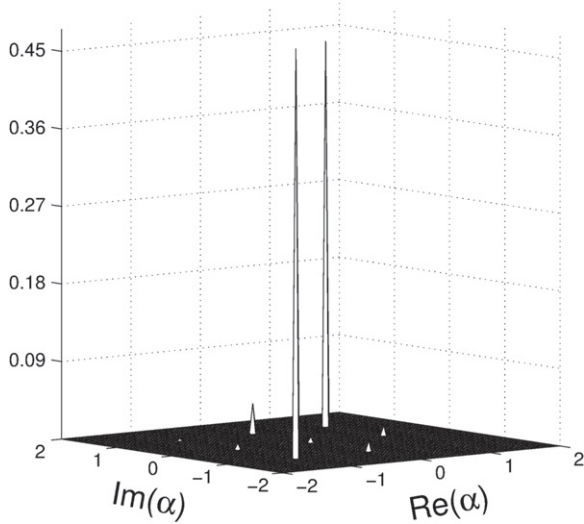


Figure 7. Pins showing the positions and the absolute values of the coefficients of the constituent coherent states for squeezed number state $|1, 0.5 \exp(i\pi/2), 0\rangle$ with $N = 9$ coherent states on a lattice.

Table 7. The misfits $\epsilon^{(\text{lattice})}$ and $\epsilon^{(\text{circle})}$ of the approximation of squeezed number states on a lattice and on a circle with 9 coherent states and the optimal distances d_{opt} and radii R_{opt} for various phase θ and a constant magnitude $r = 0.5$ of squeezing and a constant photon number $n = 1$.

State	$\epsilon^{(\text{lattice})}$	d_{opt}	$\epsilon^{(\text{circle})}$	R_{opt}
$ 1, 0.5, 0\rangle$	0.0017	1.44	0.0019	1.53
$ 1, 0.5 \exp(i\pi/12), 0\rangle$	0.0069	1.34	0.0034	1.53
$ 1, 0.5 \exp(i\pi/6), 0\rangle$	0.008	1.15	0.0012	1.63
$ 1, 0.5 \exp(i\pi/3), 0\rangle$	0.0065	1.14	0.004	1.49
$ 1, 0.5 \exp(i5\pi/12), 0\rangle$	0.0025	1.13	0.0049	1.44
$ 1, 0.5 \exp(i\pi/2), 0\rangle$	0.0033	1.04	0.0042	1.39
$ 1, 0.5 \exp(i7\pi/12), 0\rangle$	0.0032	1.07	0.0063	1.43
$ 1, 0.5 \exp(i2\pi/3), 0\rangle$	0.006	1.13	0.0044	1.42
$ 1, 0.5 \exp(i5\pi/6), 0\rangle$	0.0074	1.1	0.005	1.51
$ 1, 0.5 \exp(i11\pi/12), 0\rangle$	0.007	1.35	0.0032	1.54
$ 1, 0.5 \exp(i\pi), 0\rangle$	0.0018	1.44	0.0029	1.48

practically perfect showing the high accuracy of the approximation.

Next we consider the approximation of squeezed number states. Table 7 presents the misfits $\epsilon^{(\text{lattice})}$ and $\epsilon^{(\text{circle})}$ of the approximations, the optimal distances d_{opt} of adjacent lattice elements, and radii R_{opt} for various phases θ and for a constant magnitude $r = 0.5$ of squeezing and a constant photon number $n = 1$. The data show that the accuracy of the approximation is generally good but really high around $\theta = 0, \theta = \pi$, and $\theta = \pi/2$ with the range of $|\Delta\theta| \lesssim \pi/12$. In this case the state is elongated in the direction of the real axis, the imaginary axis, and the diagonal of the lattice, respectively. In these cases the approximation on a lattice is better than the one on a circle. Though there is no general rule in the change of the optimal distance d_{opt} and radius R_{opt} , one can observe that the optimal distances are smaller for the states with phases of squeezing around $\pi/2$ ($\theta = \pi/2 \pm \pi/6$). In

Table 8. The misfits $\epsilon^{(\text{lattice})}$ and $\epsilon^{(\text{circle})}$ of the approximation of amplitude-squeezed states on a lattice and on a circle with 9 coherent states and the optimal distances d_{opt} and radii R_{opt} for various values of u and fixed parameters $\delta = 4$ and $R = 2$.

State	$\epsilon^{(\text{lattice})}$	d_{opt}	$\epsilon^{(\text{circle})}$	R_{opt}
$ \Psi_A(0.5, 4, 2)\rangle$	0.003	0.9	5×10^{-5}	1.65
$ \Psi_A(1, 4, 2)\rangle$	0.0048	0.98	2×10^{-4}	1.74
$ \Psi_A(1.5, 4, 2)\rangle$	0.0066	1.15	4×10^{-4}	1.79
$ \Psi_A(1.6, 4, 2)\rangle$	0.0093	1.16	2.2×10^{-4}	1.72
$ \Psi_A(1.7, 4, 2)\rangle$	0.0058	1.16	9×10^{-4}	1.87
$ \Psi_A(1.9, 4, 2)\rangle$	0.0142	1.51	0.01	1.84
$ \Psi_A(2, 4, 2)\rangle$	0.0105	1.41	0.012	1.86
$ \Psi_A(3, 4, 2)\rangle$	0.011	1.9	0.0013	1.78
$ \Psi_A(5, 4, 2)\rangle$	0.0031	1.95	7×10^{-4}	1.93
$ \Psi_A(10, 4, 2)\rangle$	3×10^{-4}	1.99	1.3×10^{-4}	1.98

figure 7 we show the pins representing the positions and the absolute values of the coefficients of the constituent coherent states for squeezed number state $|1, 0.5 \exp(i\pi/2), 0\rangle$ with $N = 9$ coherent states on a lattice.

Finally, let us consider amplitude-squeezed states

$$|\Psi_A(u, \delta, R)\rangle = \mathcal{N} \int_{-\pi}^{\pi} \exp\left(-\frac{1}{2}u^2\phi^2 - i\delta\phi\right) |Re^{i\phi}\rangle d\phi \quad (27)$$

defined by Gaussian continuous coherent-state superpositions in phase space. $\mathcal{N}$ is a normalization constant.

For calculational purposes it is convenient to use the number-state representation of the states:

$$|\Psi_A(u, \delta, R)\rangle = \sum_{n=0}^{\infty} C_n |n\rangle, \quad (28)$$

where

$$C_n = \langle n | \Psi_A(u, \delta, R) \rangle = \mathcal{N} \frac{2\pi R^n}{u\sqrt{n!}} \exp\left[-\frac{(\delta - u)^2}{2u^2}\right]. \quad (29)$$

These states tend to the coherent state R in the limit $u \rightarrow \infty$ and yield the photon number state δ when δ is a non-negative integer in the limit $u \rightarrow 0$, while $\delta = R^2$, the mean values of the photon number for these limiting states are equal.

In table 8 the misfits $\epsilon^{(\text{lattice})}$ and $\epsilon^{(\text{circle})}$ of the approximations of various amplitude-squeezed states on a lattice and on a circle and the optimal parameters d_{opt} and R_{opt} are presented. The table shows that the accuracy of the approximation on both formations is high for all the considered states. The accuracy on a circle is always higher than on a lattice except for the state $|\Psi_A(2, 4, 2)\rangle$. The optimal distance d_{opt} of the lattice is growing with increasing values of the parameter u .

In figure 8 we present the pins showing the positions and the absolute values of the coefficients of the $N = 9$ constituent coherent states in the coherent-state superposition approximating the amplitude-squeezed state $|\Psi_A(1, 4, 2)\rangle$ on a

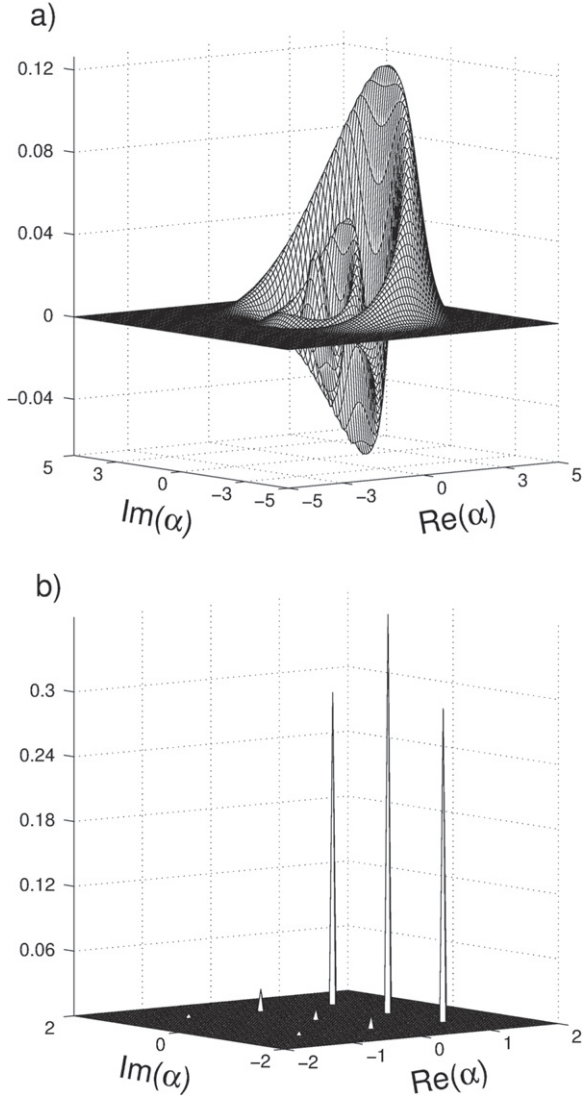


Figure 8. Amplitude-squeezed state $|\Psi_A(1, 4, 2)\rangle$ approximated by $N = 9$ coherent states on a lattice (a) Wigner function $W(\alpha)$ of the approximating coherent-state superposition, (b) pins showing the positions and the absolute values of the coefficients of the constituent coherent states.

lattice and the Wigner function of the approximating state. In this case the Wigner function is also practically perfect.

From these results one can conclude that lattice coherent-state superpositions are universal for approximating nonclassical states of harmonic oscillator systems. The accuracy of the approximation is high for several quantum states. Higher accuracy can be obtained by circle superpositions only for quantum states having circular symmetry in the Wigner function [13].

Finally, we compare the presented quantum state engineering protocol with the ones based on the number-state expansion of the target state in traveling-wave optics [8–10], with a special attention to the aspect of experimental realizability. An experimental scheme in this area generally contains beam splitters and photon number or homodyne detectors. All the schemes are probabilistic, the target state is

prepared conditionally depending on the results of the measurements. The fidelity of the prepared state depends on the perfectness of the input states of the scheme, the detector efficiencies, and the amount of unavoidable photon losses in the system. Obviously, the effects of imperfections and losses are generally proportional to the number of construction elements, e.g. the number of detectors in the experimental setup, or the number of special input states.

The experimental scheme in [45] which has inspired the analysis of quantum state engineering via coherent-state superpositions on a lattice presented in this section contains only three beam splitters and three homodyne measurements independently of the quantum state to be engineered. As the number of elements in the experimental schemes for quantum state engineering based on number-state expansion is proportional to the number of terms considered in the expansion, the presented coherent-state superposition based method can be generally advantageous for quantum states containing several terms in their number-state expansion.

We note that the effects of experimental imperfections have been analyzed in the literature for various schemes developed for generating Schrödinger cat states in traveling-wave optics [25, 32, 64–66]. Such schemes have similar structure to those capable of producing general states. For example, the experimental scheme for generating a general quantum state with a finite number-state expansion using photon subtraction can be considered as a generalization of the scheme developed for producing Schrödinger cat states by the same process. These analyses show that the change in the fidelity characterizing the performance of the engineering due to experimental imperfections strongly depends on the specifics of the given scheme and even on the properties of the target state. Generally, experimental schemes with homodyne detections might have some advantage because this type of measurement can be performed with much higher efficiency than photon counting measurements [65, 66]. This property can yield an additional benefit for coherent-state superposition based quantum engineering protocols because using homodyne detections can be plausible in these schemes and it also holds for the experimental scheme in [45]. We hope that the presented results will motivate further research on developing experimental schemes capable of producing coherent-state superpositions approximating nonclassical states in traveling-wave optics.

5. Conclusions

We have considered the optimal approximation of certain quantum states of a harmonic oscillator with the superposition of a finite number of coherent states on an ellipse and on a lattice in phase space. All the coherent-state superpositions we have considered are feasible in current experiments. We have optimized numerically the parameters of the chosen geometry and the coefficients of the coherent states via a genetic algorithm, in order to obtain the best feasible approximation.

First, we placed the coherent states equidistantly in their phase on an ellipse. As the ellipse fits to the shape of the Wigner function of certain states, we have expected that for such states we can obtain an approximation superior in quality to the one based on the one-dimensional coherent-state representation of the state, that is, constructed from states along a line in phase space. We have found this intuition to be appropriate. In particular, for squeezed number states in a certain parameter range, the elliptical approximation outperforms the one on a line.

Next we have considered a 3×3 equidistant lattice around the origin of phase space. These are the kinds of superpositions which seem to be feasible in traveling-wave optical experiments. A relatively large set of states, including special number-state superpositions, appear to be better approximated in this geometry than with a superposition on a circle.

Acknowledgments

We are sincerely grateful to Margarita A Man'ko and Vladimir I Man'ko for their long-standing collaboration with our group. We are grateful for the support of the Hungarian Scientific Research Fund OTKA (Contract No. K83858). We thank M Koniorczyk for useful discussions.

References

- [1] Dodonov V V 2002 *J. Opt. B: Quantum Semiclass. Opt.* **4** 1–33
- [2] Dodonov V V, Man'ko M A, Man'ko V I and Vourdas A 2007 *J. Russ. Laser Res.* **28** 404
- [3] Dell'Anno F, De Siena S and Illuminati F 2006 *Phys. Rep.* **428** 53–168
- [4] Vogel K, Akulin V M and Schleich W P 1993 *Phys. Rev. Lett.* **71** 1816
- [5] Parkins A S, Marte P, Zoller P and Kimble H J 1993 *Phys. Rev. Lett.* **71** 3095
- [6] Parkins A S, Marte P, Zoller P, Carnal O and Kimble H J 1995 *Phys. Rev. A* **51** 1578
- [7] Law C K and Eberly J H 1996 *Phys. Rev. Lett.* **76** 1055
- [8] Dakna M, Clausen J, Knöll L and Welsch D-G 1999 *Phys. Rev. A* **59** 1658
- [9] Villas-Boas C J, Guimarães Y, Moussa M H Y and Baseia B 2001 *Phys. Rev. A* **63** 055801
- [10] Fiurášek J, García-Patrón R and Cerf N J 2005 *Phys. Rev. A* **72** 033822
- [11] Janszky J, Domokos P and Adam P 1993 *Phys. Rev. A* **48** 2213
- [12] Janszky J, Domokos P, Szabó S and Adam P 1995 *Phys. Rev. A* **51** 4191
- [13] Szabo S, Adam P, Janszky J and Domokos P 1996 *Phys. Rev. A* **53** 2698
- [14] Dodonov V V, Malkin I A and Man'ko V I 1974 *Physica* **72** 597
- [15] Bužek V, Vidiella-Barranco A and Knight P L 1992 *Phys. Rev. A* **45** 6570
- [16] Dodonov V V, Yu K S and Man'ko V I 1995 *Phys. Lett. A* **199** 123
- [17] Dodonov V V, Man'ko V I and Nikonov D E 1995 *Phys. Rev. A* **51** 3328
- [18] Dodonov V V and Man'ko V I 2003 *Theory of Nonclassical States of Light* (London, Boca Raton, FL: Taylor and Francis, CRC Press) Ch 1, 4
- [19] Brune M, Hagley E, Dreyer J, Mai'etre X, Maali A, Wunderlich C, Raimond J M and Haroche S 1996 *Phys. Rev. Lett.* **77** 4887
- [20] Auffeves A, Maioli P, Meunier T, Gleyzes S, Nogues G, Brune M, Raimond J M and Haroche S 2003 *Phys. Rev. Lett.* **91** 230405
- [21] Jeong H, Lance A M, Grosse N B, Symul T, Lam P K and Ralph T C 2006 *Phys. Rev. A* **74** 033813
- [22] Neegaard-Nielsen J S, Nielsen B M, Hettich C, Molmer K and Polzik E S 2006 *Phys. Rev. Lett.* **97** 083604
- [23] Ourjoumtsev A, Jeong H, Tualle-Brouiri R and Grangier P 2007 *Nature* **448** 784
- [24] Nielsen A E B and Molmer K 2007 *Phys. Rev. A* **76** 043840
- [25] Glancy S and Vasconcelos H M 2008 *J. Opt. Soc. Am. B* **25** 712
- [26] Takahashi H, Wakui K, Suzuki S, Takeoka M, Hayasaka K, Furusawa A and Sasaki M 2008 *Phys. Rev. Lett.* **101** 233605
- [27] He B, Nadeem M and Bergou J A 2009 *Phys. Rev. A* **79** 035802
- [28] Adam P, Kiss T, Mechler M and Darazs Z 2010 *Phys. Scr. T* **140** 014011
- [29] Adam P, Kiss T, Darazs Z and Jex I 2011 *Phys. Scr. T* **143** 014002
- [30] Lee C-W, Lee J, Nha H and Jeong H 2012 *Phys. Rev. A* **85** 063815
- [31] Wang S, Yuan H and Xu X 2013 *Eur. Phys. J. D* **67** 102
- [32] Laghaout A, Neegaard-Nielsen J S, Rigas I, Kragh C, Tjpsmark A and Andersen U L 2013 *Phys. Rev. A* **87** 043826
- [33] Brune M, Haroche S, Raimond J M, Davidovich L and Zagury N 1992 *Phys. Rev. A* **45** 5193
- [34] Guo G C and Zheng S B 1996 *Phys. Lett. A* **223** 332
- [35] Wang Z J 2007 *Phys. Rev. A* **76** 043403
- [36] Domokos P, Janszky J and Adam P 1994 *Phys. Rev. A* **50** 3340
- [37] Lutterbach L G and Davidovich L 2000 *Phys. Rev. A* **61** 023813
- [38] Maia L P A, Baseia B, Avelar A T and Malbouisson J M C 2004 *J. Opt. B: Quantum Semiclass. Opt.* **6** 351
- [39] Pathak P K and Agarwal G S 2005 *Phys. Rev. A* **71** 043823
- [40] Avelar A T, Da Rocha Filho T M, Losano L and Baseia B 2005 *Phys. Lett. A* **340** 74
- [41] Zheng S B 2007 *Opt. Commun.* **273** 460
- [42] Yang Z B, Zhang B and Zheng S B 2010 *Opt. Commun.* **283** 2872
- [43] Moya-Cessa H, Wallentowitz S and Vogel W 1999 *Phys. Rev. A* **59** 2920
- [44] José W D and Mizrahi S S 2000 *J. Opt. B: Quantum Semiclass. Opt.* **2** 306
- [45] Molnar E, Varga A, Mogyorosi G and Adam P 2014 *21st Central European Workshop on Quantum Optics (Brussels)* p 213 (book of abstracts)
- [46] Stergioulas L K, Vassiliadis V S and Vourdas A 1999 *J. Phys. A: Math. Gen.* **32** 3169
- [47] Stergioulas L K and Vourdas A 1999 *J. Mod. Opt.* **45** 1155
- [48] Arevalo Aguilar L M and Moya-Cessa H 2004 *Phys. Scr.* **70** 14
- [49] Wang Y, Liao Q, Liu Z, Wang J and Liu S 2011 *Opt. Commun.* **284** 282
- [50] Domokos P, Adam P and Janszky J 1994 *Phys. Rev. A* **50** 4293
- [51] Janszky J and Vinogradov A V 1990 *Phys. Rev. Lett.* **64** 2771
- [52] Adam P, Janszky J and Vinogradov A V 1990 *Opt. Commun.* **80** 155
- [53] Adam P and Janszky J 1990 *Phys. Lett. A* **149** 67
- [54] Adam P, Janszky J and Vinogradov A V 1991 *Phys. Lett. A* **160** 506
- [55] Bužek V and Knight P L 1991 *Opt. Commun.* **81** 331
- [56] Adam P, Foldesi I and Janszky J 1994 *Phys. Rev. A* **49** 1281

- [57] Szabo S, Adam P and Janszky J 2000 *Laser Phys.* **10** 139
- [58] Goldberg D E 1989 *Genetic Algorithms in Search, Optimization and Machine Learning* (Boston, MA: Addison-Wesley)
- [59] Ralph T C, White A G, Munro W J and Milburn G J 2001 *Phys. Rev. A* **65** 012314
- [60] Adesso G, Dell'Anno F, De Siena S, Illuminati F and Souza L A M 2009 *Phys. Rev. A* **79** 040305 R
- [61] Marek P and Fiurášek J 2009 *Phys. Rev. A* **79** 062321
- [62] Pegg D T, Phillips L S and Barnett S M 1998 *Phys. Rev. Lett.* **81** 1604
- [63] Lee S-Y and Nha H 2010 *Phys. Rev. A* **82** 053812
- [64] Suzuki S, Tsujino K, Kannari F and Sasaki M 2006 *Opt. Commun.* **259** 758
- [65] Yeong H, Kim M S, Ralph T C and Ham B S 2004 *Phys. Rev. A* **70** 061801 R
- [66] Yeong H, Lund A P and Ralph T C 2005 *Phys. Rev. A* **72** 013801